

Ultrafast X-ray and Molecular Diffraction

*Philip Bucksbaum
Stanford PULSE Institute
SLAC National Accelerator Laboratory,
Stanford University
Menlo Park, California, USA*

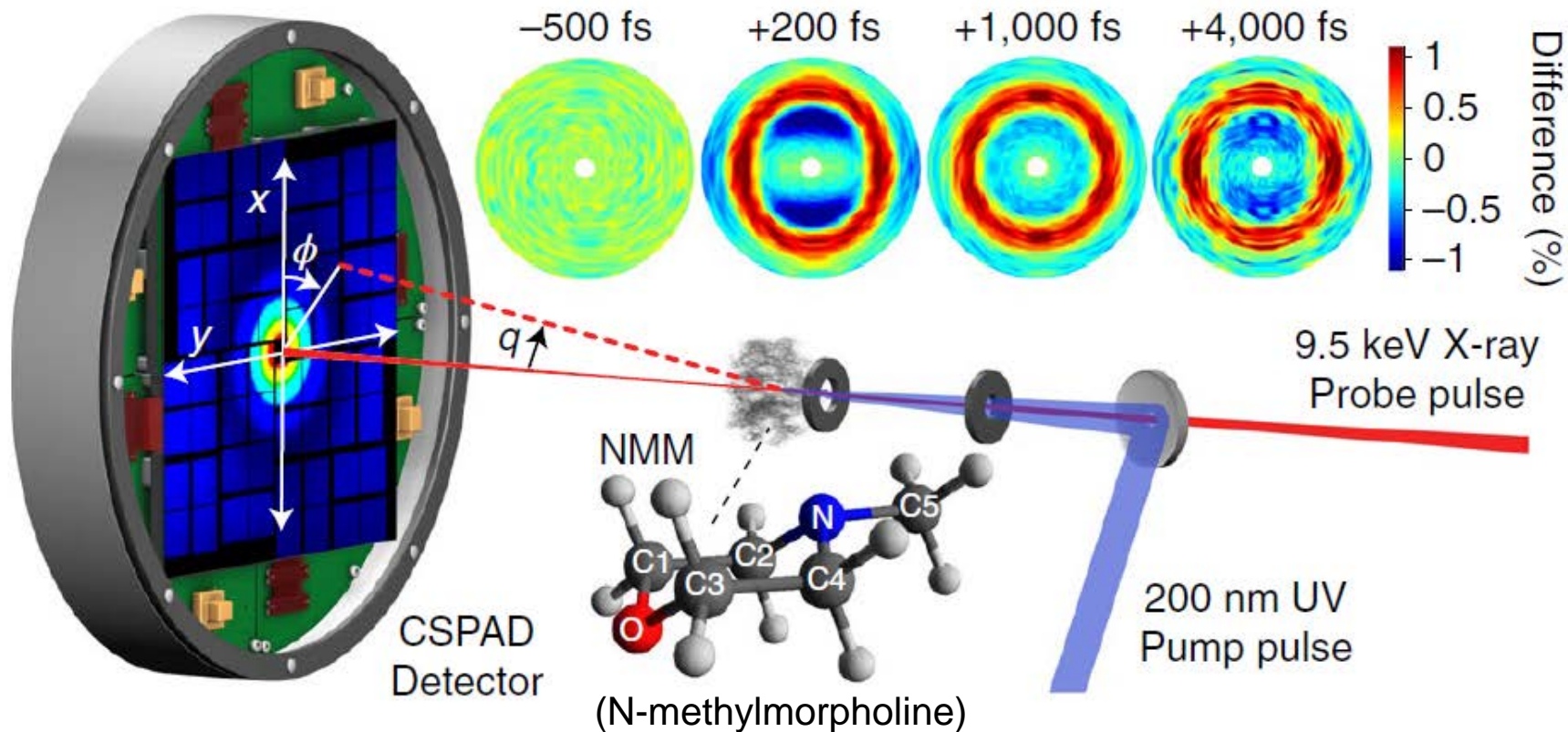
**Matthew Ware
Mike Glownia
Adi Natan
James Cryan
Ian Gabalski
PHB**

Stanford

**PULSE
INSTITUTE**

Parts of this work were supported by the U.S. Department of Energy, Office of Science, Chemical Sciences, Geosciences, and Biosciences Division

Several groups are now making femtosecond movies of intramolecular dynamics:



B. Stankus, H. Yong, N. Zotev, J. M. Ruddock, D. Bellshaw, T. J. Lane, M. Liang, S. Boutet, S. Carbajo, J. S. Robinson, W. Du, N. Goff, Y. Chang, J. E. Koglin, M. P. Minitti, A. Kirrander, and P. M. Weber, *Nat. Chem.* 11, 716 (2019).

Hard X-ray Elastic Scattering

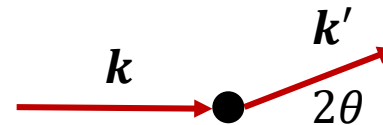
Interaction Hamiltonian:
Elastic Nonresonant scattering

$$H_I = \frac{e}{m} \mathbf{p} \cdot \mathbf{A} + \frac{e}{m} \mathbf{S} \cdot \mathbf{B} + \frac{e^2}{2m} \mathbf{A}^2$$

Scattering theory: Lowest order
Single scattering (Born) approximation

$$\text{Field Amplitude} = e^{ikz} + f(\theta) \frac{e^{ikr}}{r}$$

$$|\mathbf{Q}| = |\mathbf{k} - \mathbf{k}'| = 2k \sin \theta$$



Scattering
matrix element:

$$f(\vec{Q}) = \epsilon \cdot \epsilon' \sum_n \langle \varphi_n | e^{i\vec{Q} \cdot \vec{r}_n} | \varphi_n \rangle$$

Scattered intensity:
(matrix element)²

$$|f(\vec{Q})|^2 = \sigma_{Th}(\Omega) \left| \sum_n \langle \varphi_n | e^{i\vec{Q} \cdot \vec{r}_n} | \varphi_n \rangle \right|^2$$

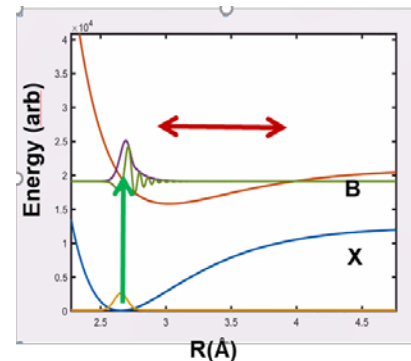
$$= \sigma_{Th}(\Omega) \sum_{i,j} \langle \psi_m(\vec{r}_i \dots \vec{r}_j) | e^{i\vec{Q} \cdot (\vec{r}_i - \vec{r}_j)} | \psi_m(\vec{r}_i \dots \vec{r}_j) \rangle$$

Scattering measures the relative position of pairs of electrons.

Approximations to simplify this further

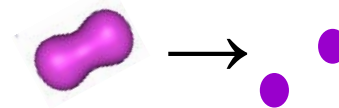
Born-Oppenheimer approximation...

Nuclei move on potential energy surfaces $\chi_n(\vec{R}, \tau)$



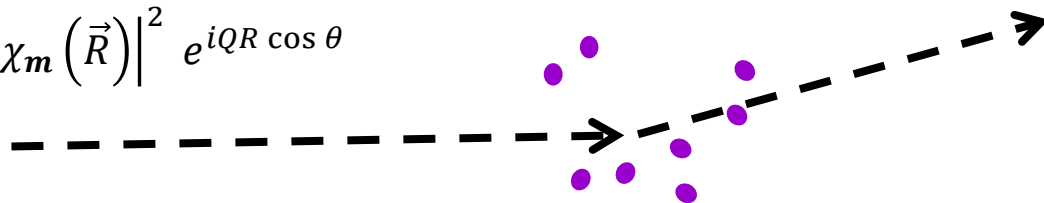
Independent atom approximation...

X-rays scatter from these atomic locations $\chi_n(\vec{R}, \tau)$



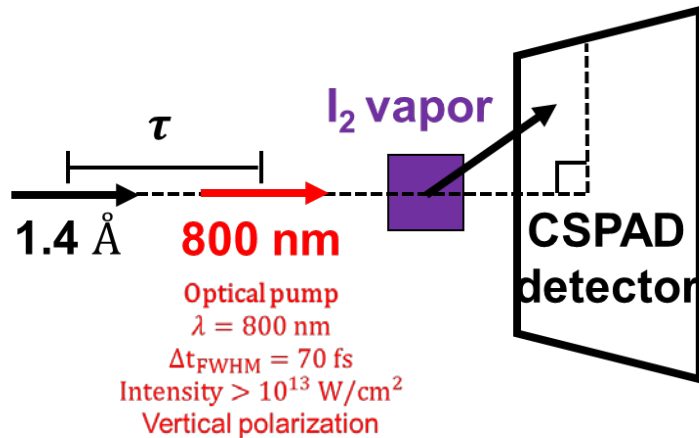
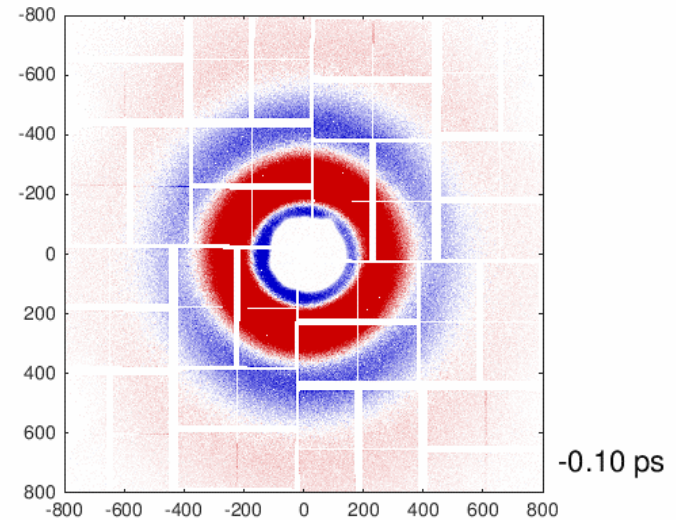
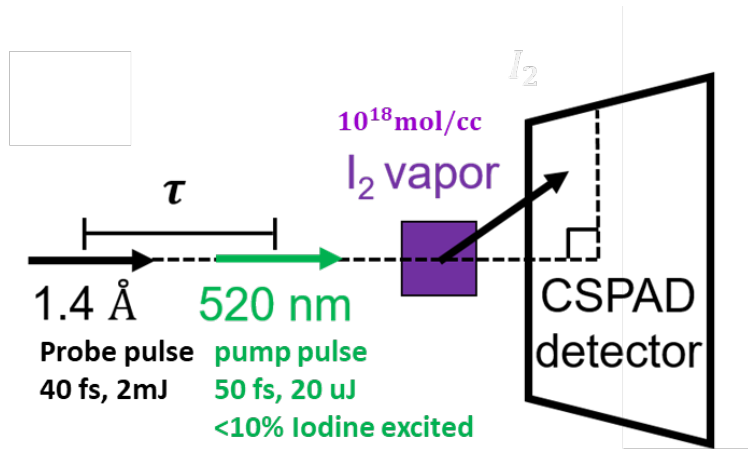
Homonuclear diatomic molecule in an eigenstate...

$$\langle \psi_m | \hat{S}(\vec{Q}) | \psi_m \rangle = 2 |f_I(Q)|^2 \int d\vec{R} |\chi_m(\vec{R})|^2 e^{iQR \cos \theta}$$

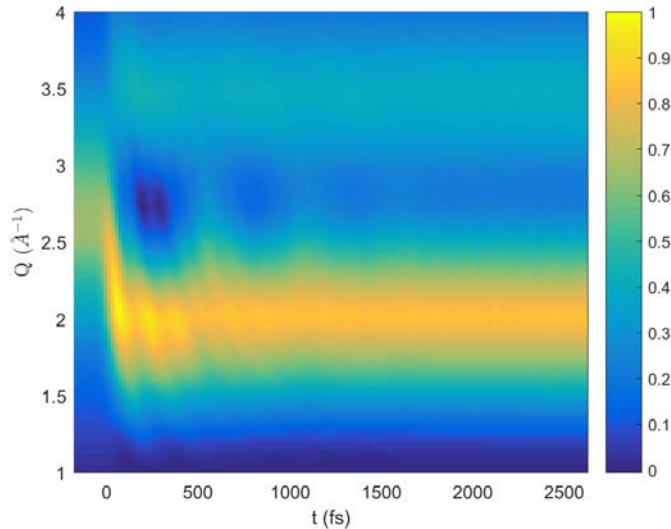


$\rho_m(r)$ is the cosine transform of $\langle \hat{S}(\vec{Q}) \rangle$.

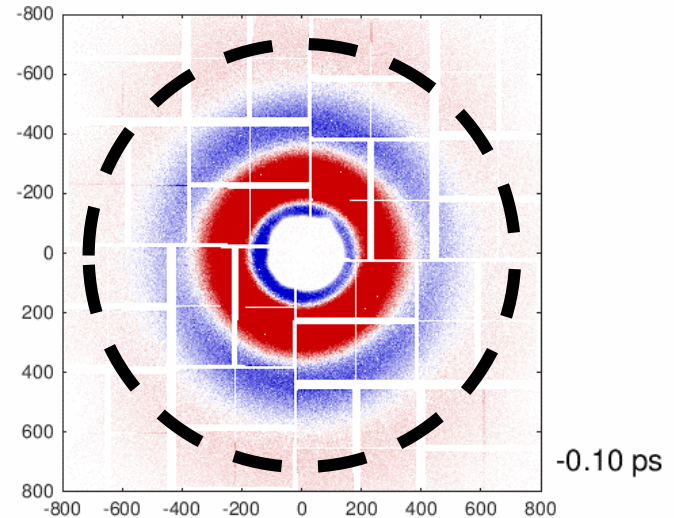
Here's the real thing:



Making movies: Angle-integrated signal



**($t < 0$) signal partly
subtracted to keep
signal non-negative
everywhere**

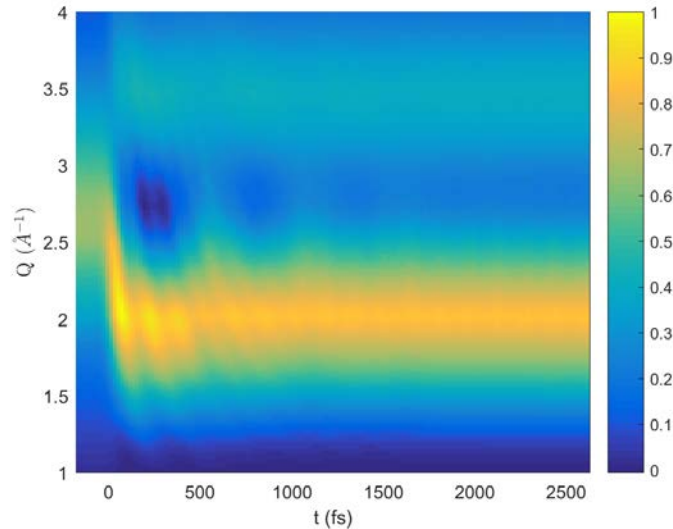


$$Q(R) = \oint Q(R, \phi) d\phi$$

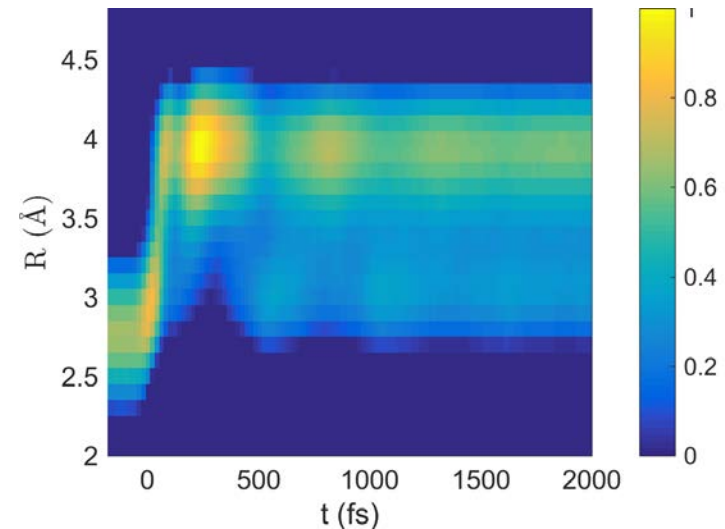
Mike Glownia, Adi Natan, James Cryan,
et al., Phys. Rev. Lett. 117, 153003
(2016); Reply to SM Comment, 119,
069302 (2017).

Natan, Ware

Making movies: Angle-integrated signal



**($t < 0$) signal partly
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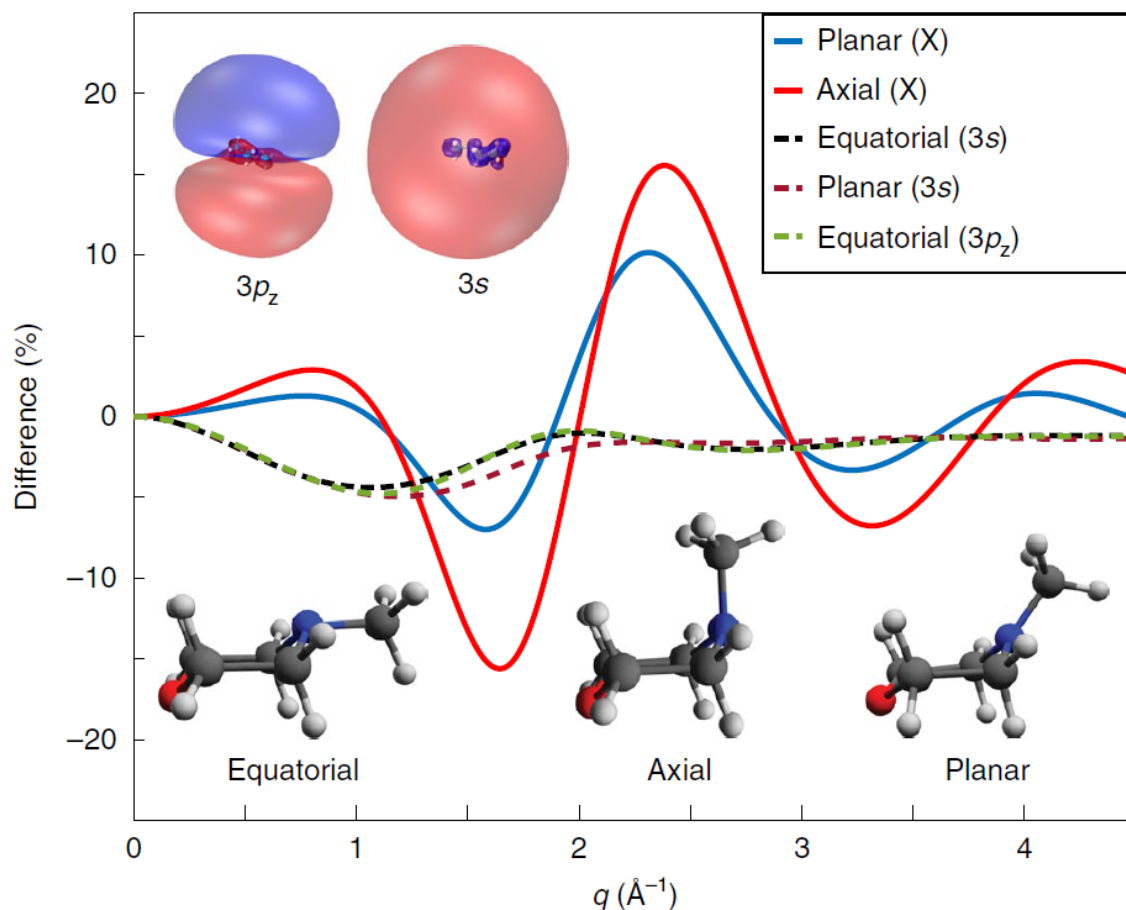


**R-space movie.
The signal at $t < 0$ is
the fraction of the X
state that was
excited at $t=0$.**

Mike Glownia, Adi Natan, James Cryan,
et al., Phys. Rev. Lett. 117, 153003
(2016); Reply to SM Comment, 119,
069302 (2017).

Natan, Ware

Polyatomic molecules are more challenging: First, there are lots of conformers:



Calculated difference scattering for this molecule (N-methylmorpholine)

Blue solid line: Planar vs Equatorial in the ground state.

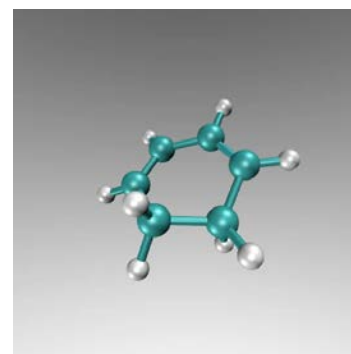
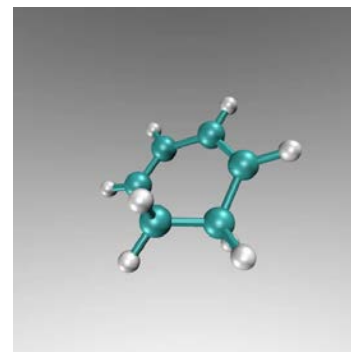
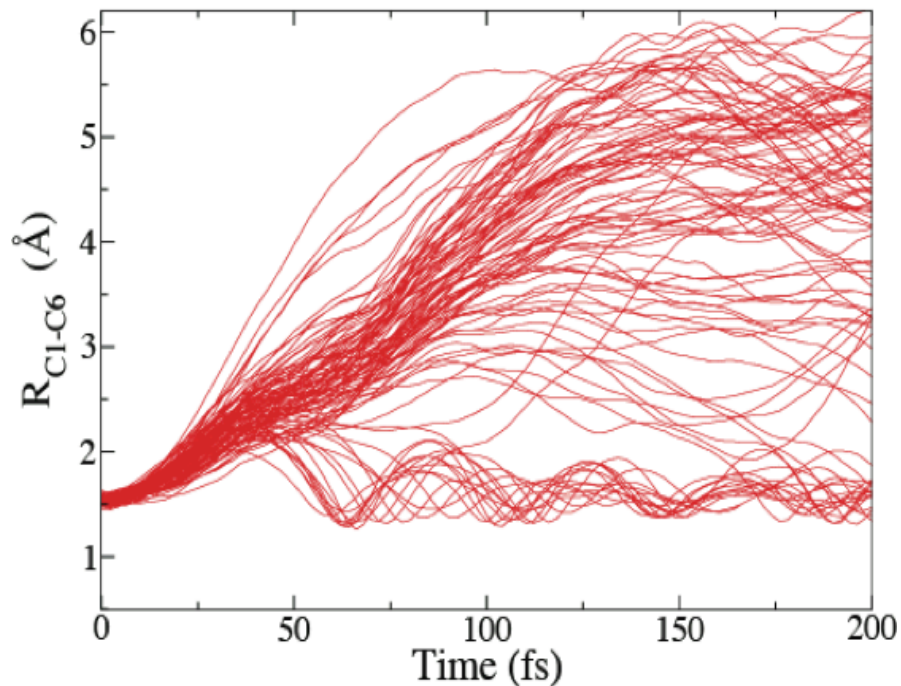
Red solid line: axial vs. equatorial in the ground state,

Dotted lines: Excitation of 3s and 3p orbitals

B. Stankus, H. Yong, N. Zotev, J. M. Ruddock, D. Bellshaw, T. J. Lane, M. Liang, S. Boutet, S. Carbajo, J. S. Robinson, W. Du, N. Goff, Y. Chang, J. E. Koglin, M. P. Minitti, A. Kirrander, and P. M. Weber, *Nat. Chem.* 11, 716 (2019).

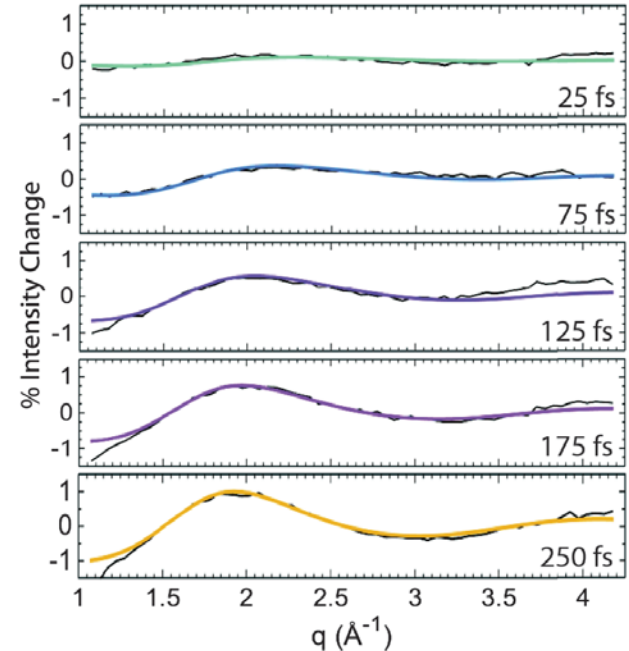
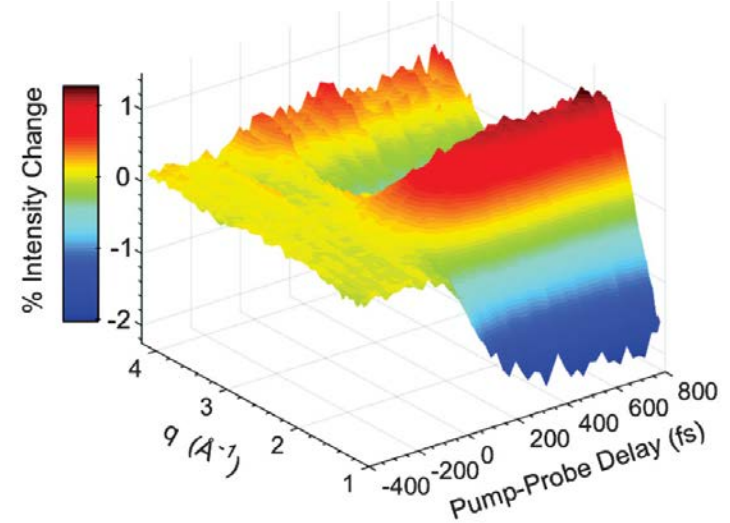
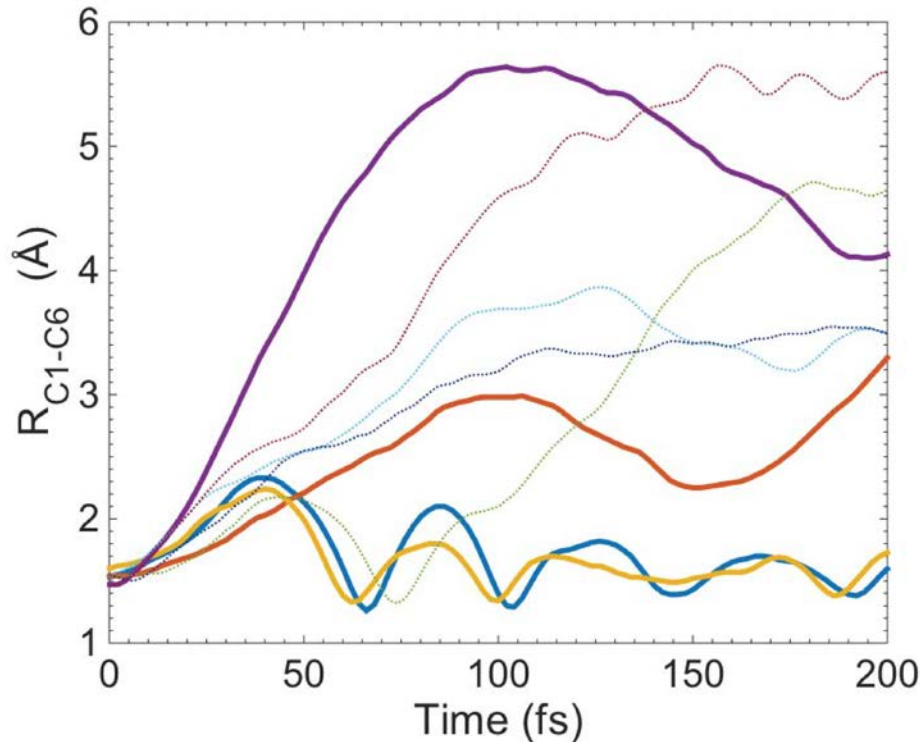
Also lots of trajectories.

Cyclohexadiene



M. P. Minitti, J. M. Budarz, A. Kirrander, J. S. Robinson, D. Ratner, T. J. Lane, D. Zhu, J. M. Glowia, M. Kozina, H. T. Lemke, M. Sikorski, Y. Feng, S. Nelson, K. Saita, B. Stankus, T. Northey, J. B. Hastings, and P. M. Weber, *Phys. Rev. Lett.* **114**, 255501 (2015).

A current approach: Find trajectories that fit the data best:

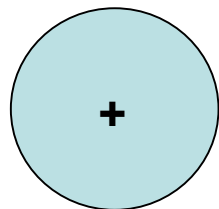
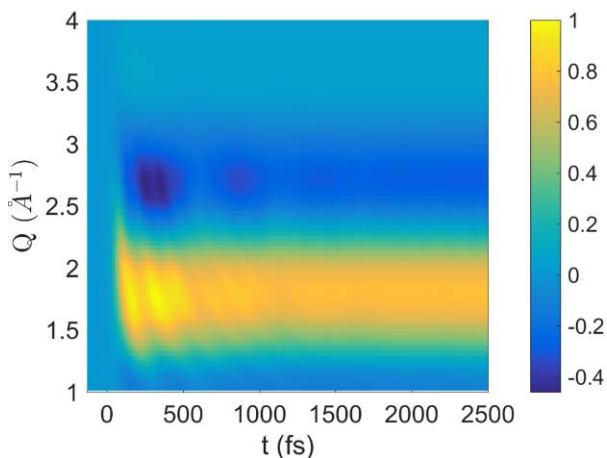


Minitti, M. P. *et al. Phys. Rev. Lett.* 114, 255501 (2015).

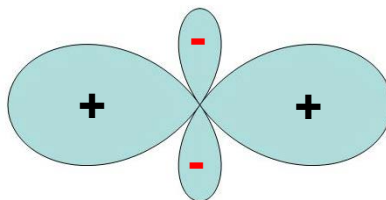
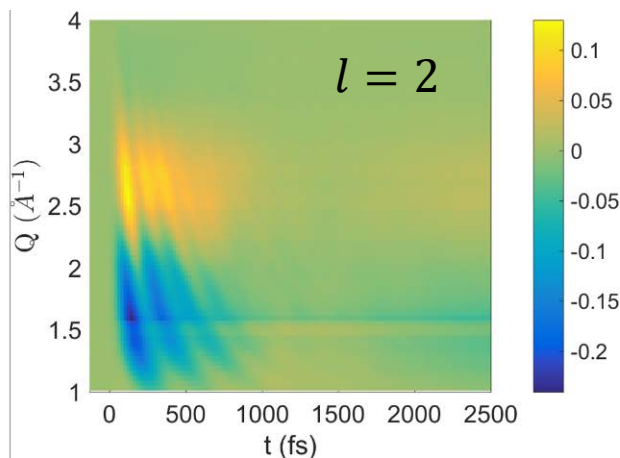
Another approach: Align the molecule to the lab frame, for example using dipole selection rules:

$$S_N(Q, \theta, \tau) = \sum_{l=0,2,\dots} \sqrt{\frac{2\pi}{2l+1}} P_l(\cos \theta) S_{N,l}(Q, \tau).$$

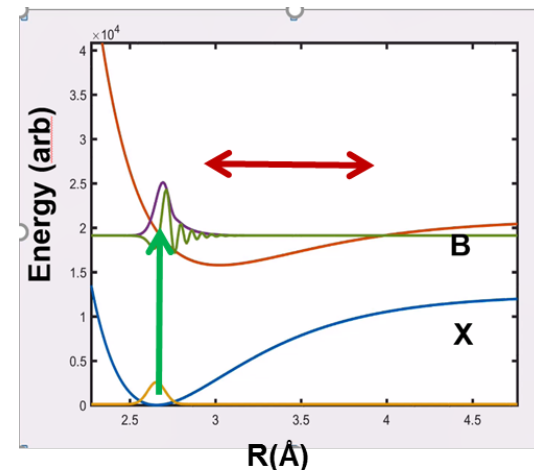
$S_0(Q, t)$



$S_2(Q, t)$



Linear polarization

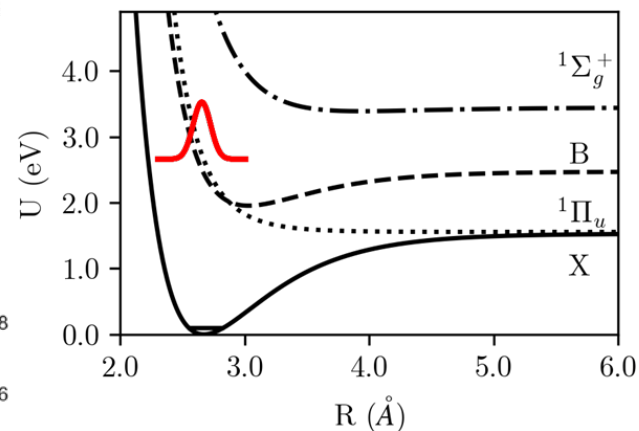
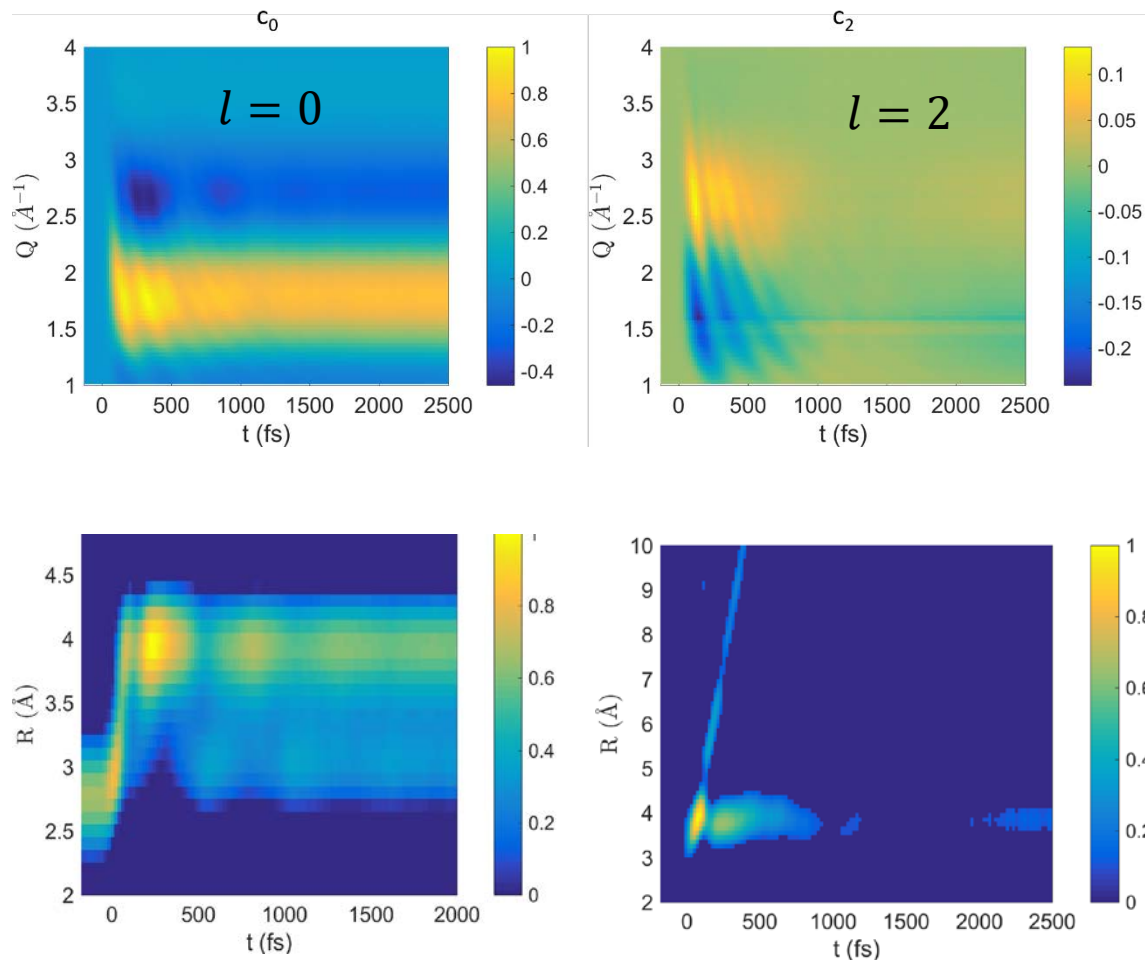


$$S_X(Q, t) \sim \text{isotropic} \sim S_0$$

$$S_B \sim \cos^2 \theta \sim S_0 + 2S_2$$

Mike Glownia, Adi Natan, James Cryan,
et al., Phys. Rev. Lett. 117, 153003
(2016); Reply to SM Comment, 119,
069302 (2017).

Reconstructions emphasize different aspects of the wave packet

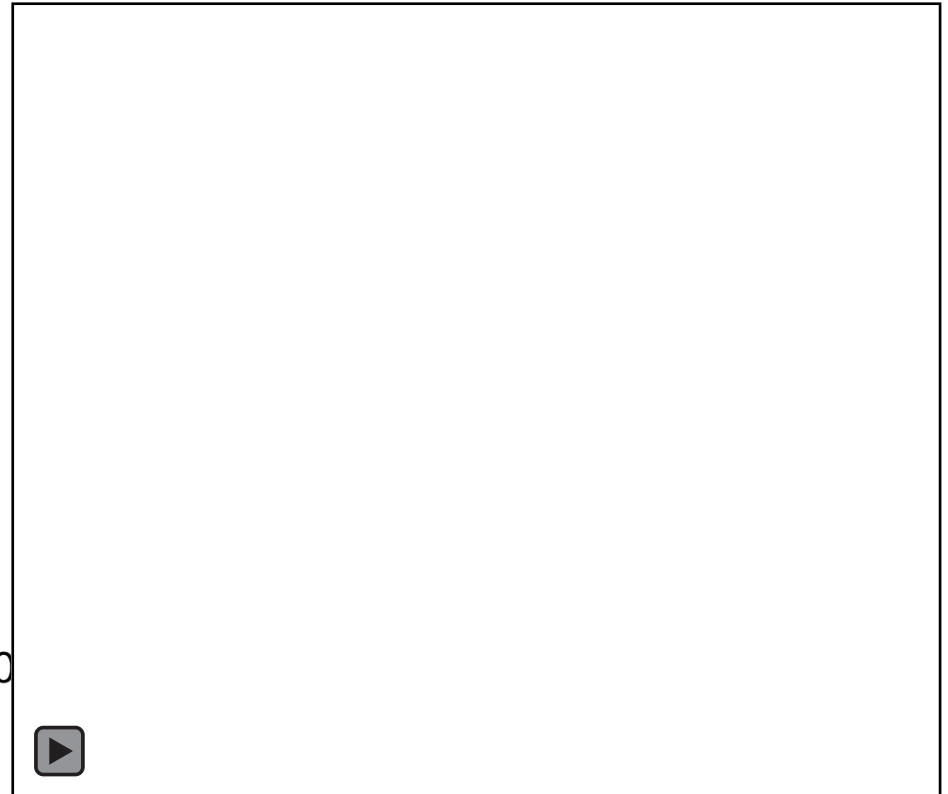
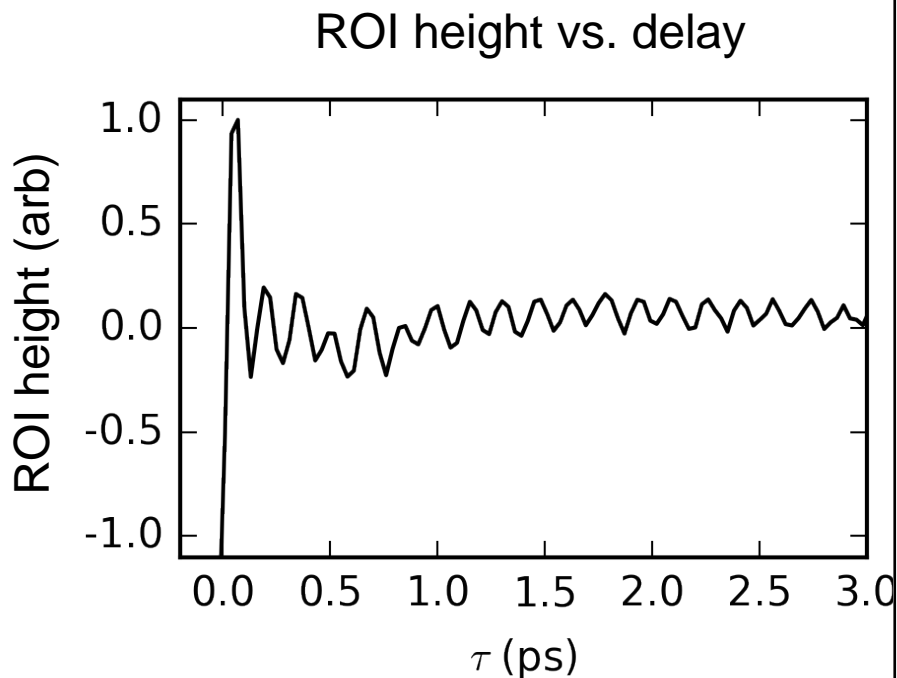


X-state subtracted

B'' dissociation

Mike Glownia, Adi Natan, James Cryan,
et al., Phys. Rev. Lett. 117, 153003
(2016); Reply to SM Comment, 119,
069302 (2017).

A third approach: Each angle has its own special time-dependence:

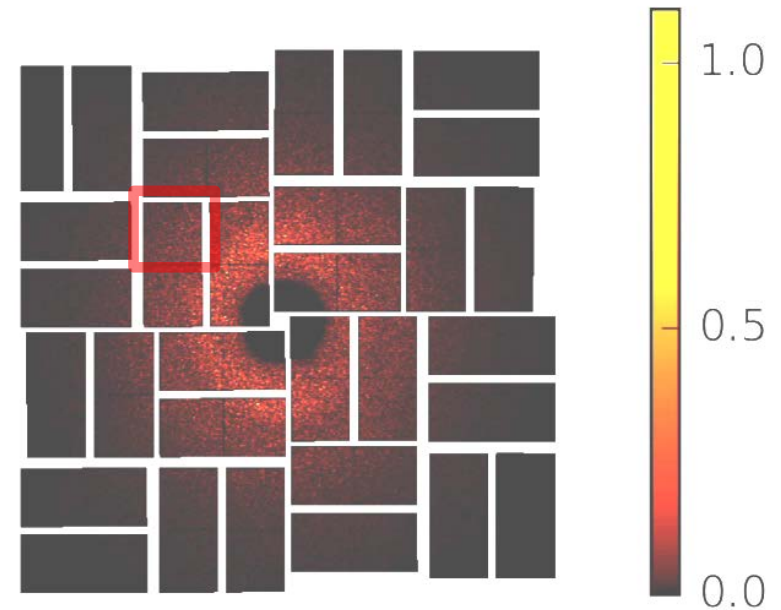
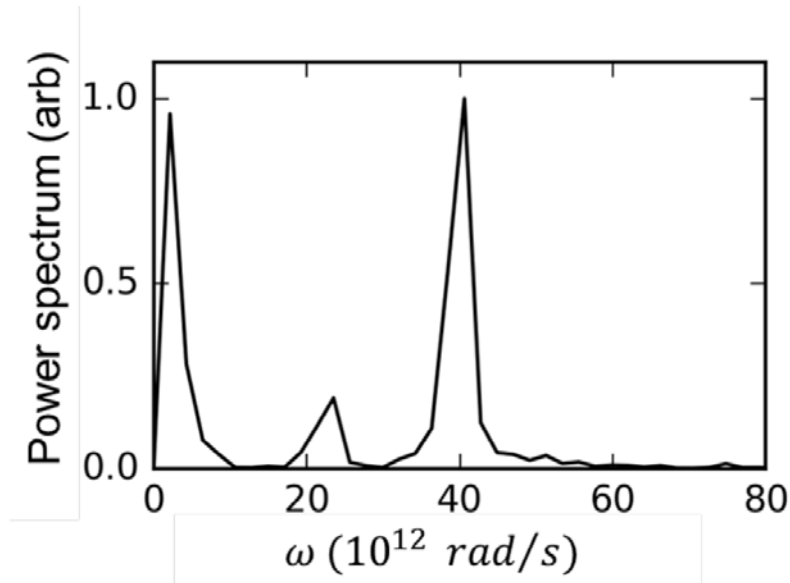


M. R. Ware, J. M. Glowina, A. Natan,
J. P. Cryan, and P. H. Bucksbaum,
Philosophical Transactions of the
Royal Society A: Mathematical,
Physical and Engineering Sciences
377, 20170477 (2019).

Matt Ware

Frequency-resolved x-ray scattering (FRXS) from molecular iodine

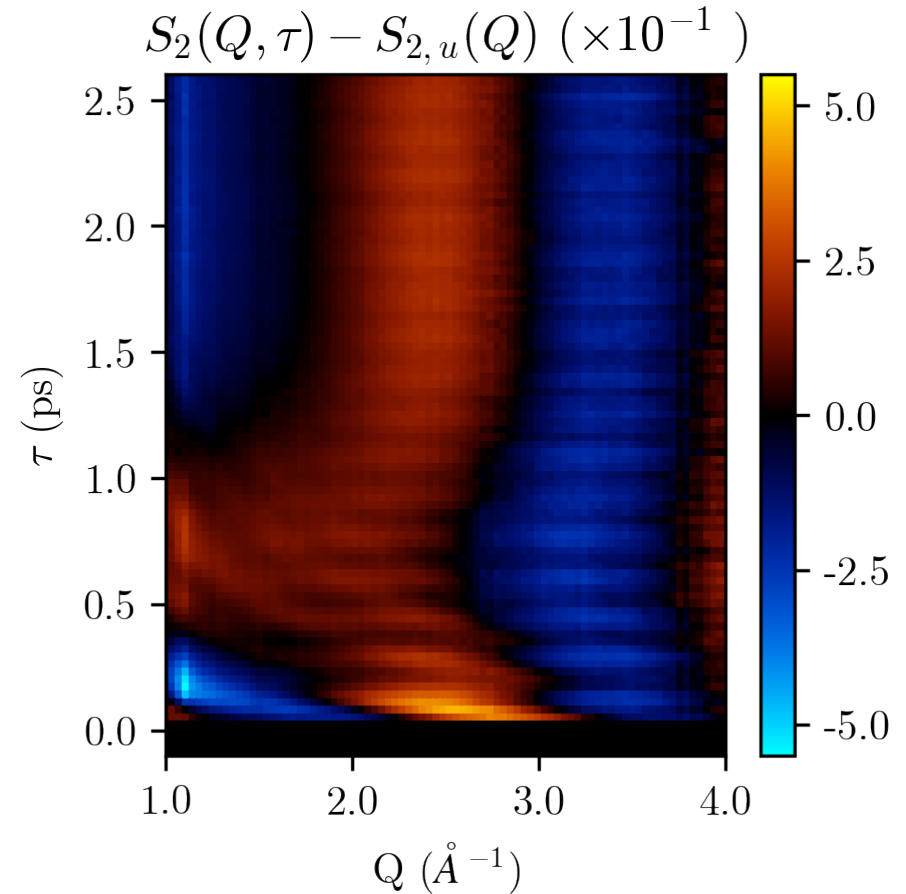
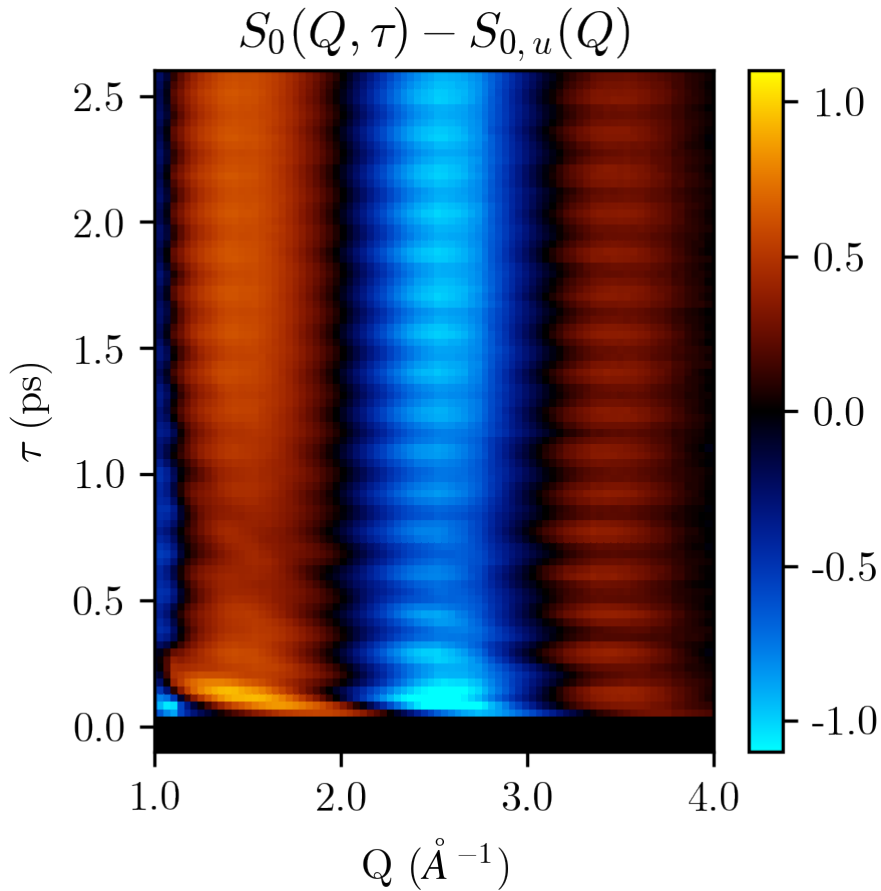
39.6 THz



$$|\tilde{S}(x, y, \omega)|^2 = \left| \int_{-\infty}^{+\infty} d\tau e^{-i\omega\tau} S(x, y, \tau) \right|^2$$
$$\Delta\omega = \frac{2\pi}{\tau_{Range}} \quad \omega_{max} = \frac{2\pi}{\Delta\tau}$$

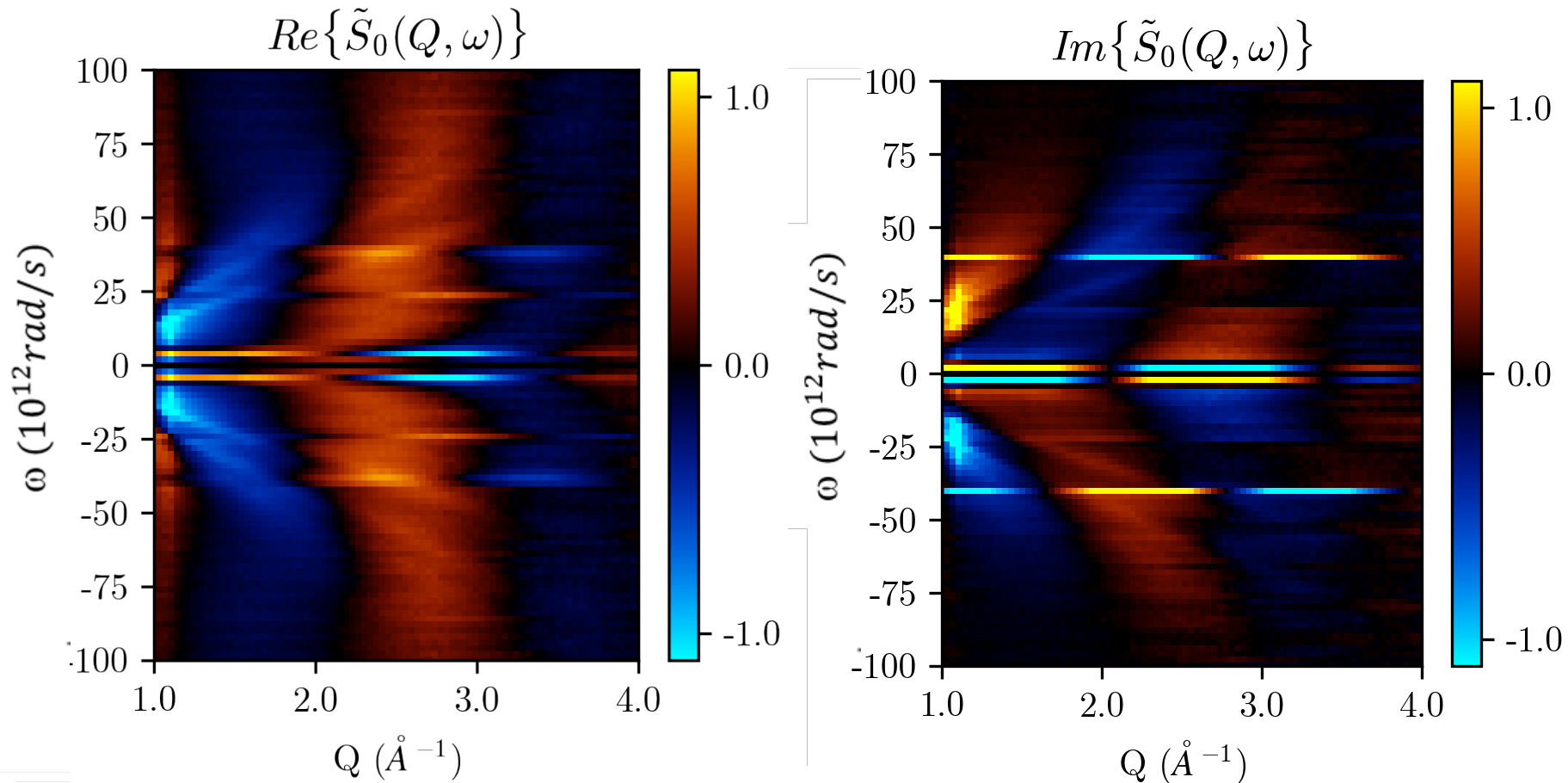
M. R. Ware, J. M. Glowina, A. Natan, J. P. Cryan, and P. H. Bucksbaum, Philosophical Transactions of the Royal Society A: Mathematical, Physical and Engineering Sciences 377, 20170477 (2019).

Apply this to Legendre decompositions.



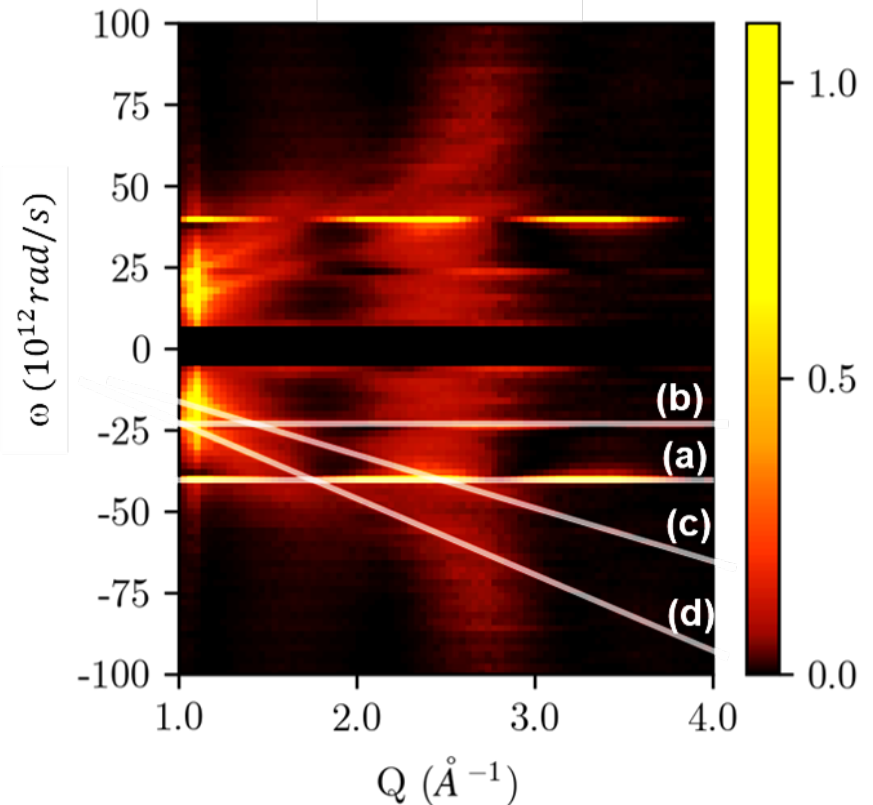
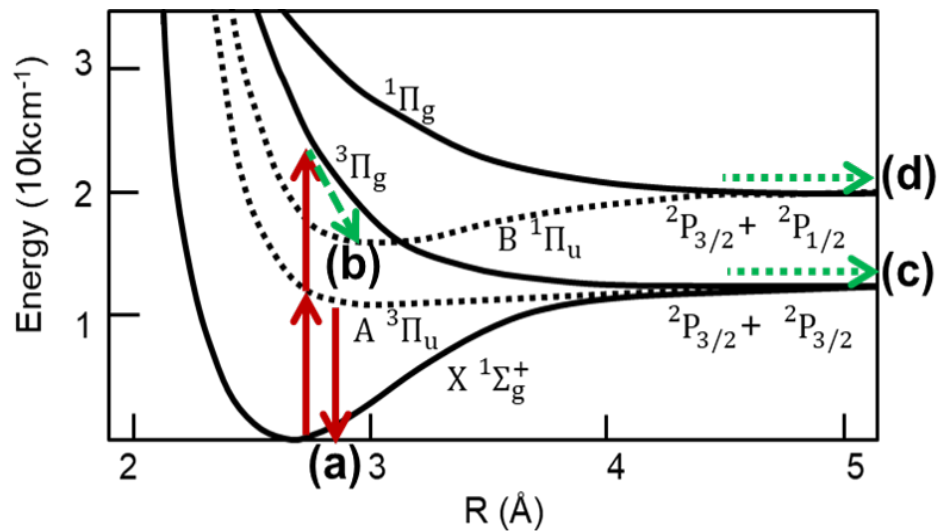
$$S_n(Q, \tau) = \int dR R^2 j_n(QR) \rho_n(R, \tau)$$

Frequency-resolved x-ray scattering spectral amplitudes



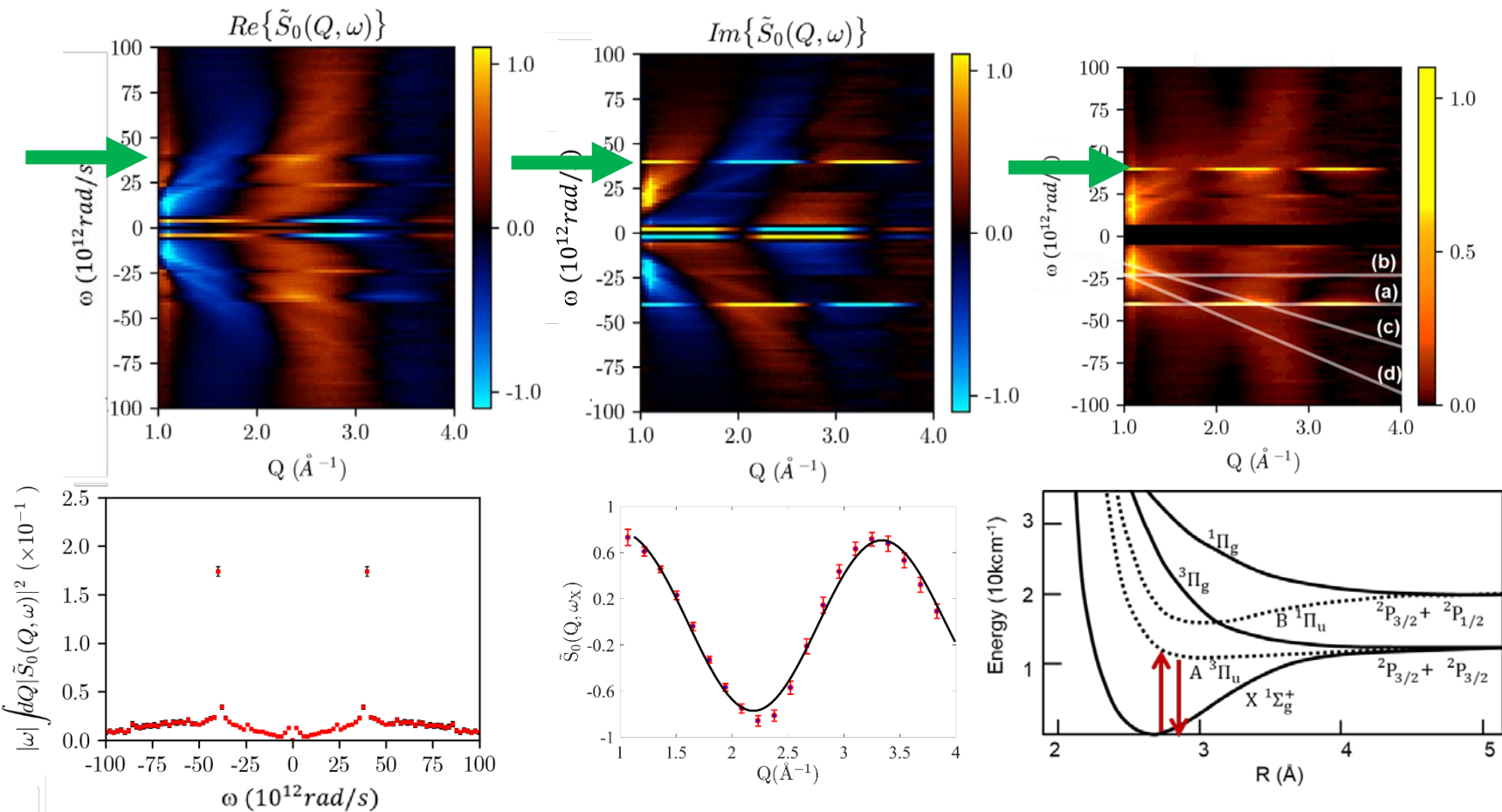
P. H. Bucksbaum, M. R.
Ware, A. Natan, J. P.
Cryan, and J. M. Glownia,
ArXiv:1911.01323
[Physics] (2019).

Power spectrum isolates different modes



P. H. Bucksbaum, M. R. Ware, A. Natan, J. P. Cryan, and J. M. Glownia, ArXiv:1911.01323 [Physics] (2019).

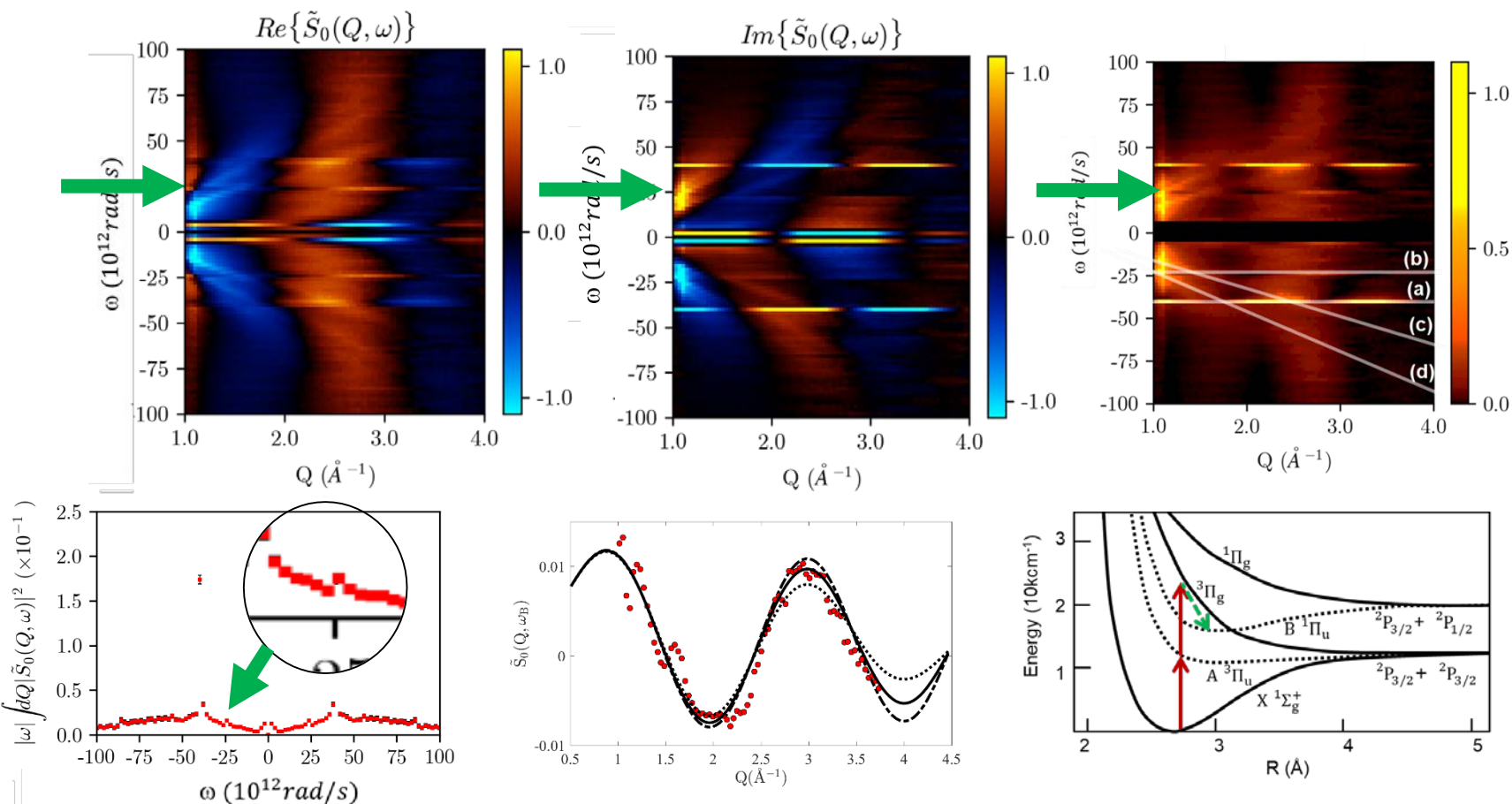
Impulsive Stimulated Raman Scattering



Center of oscillation	$2.79 \pm 0.07 \text{ \AA}$
Oscillation frequency ω	$40.3 \pm 1.0 \times 10^{12} \text{ rad/s}$
Oscillation amplitude	$0.164 \pm 0.009 \text{ \AA}$
Oscillation phase δ	290°

P. H. Bucksbaum, M. R. Ware,
A. Natan, J. P. Cryan, and J.
M. Glowia, ArXiv:1911.01323
[Physics] (2019).

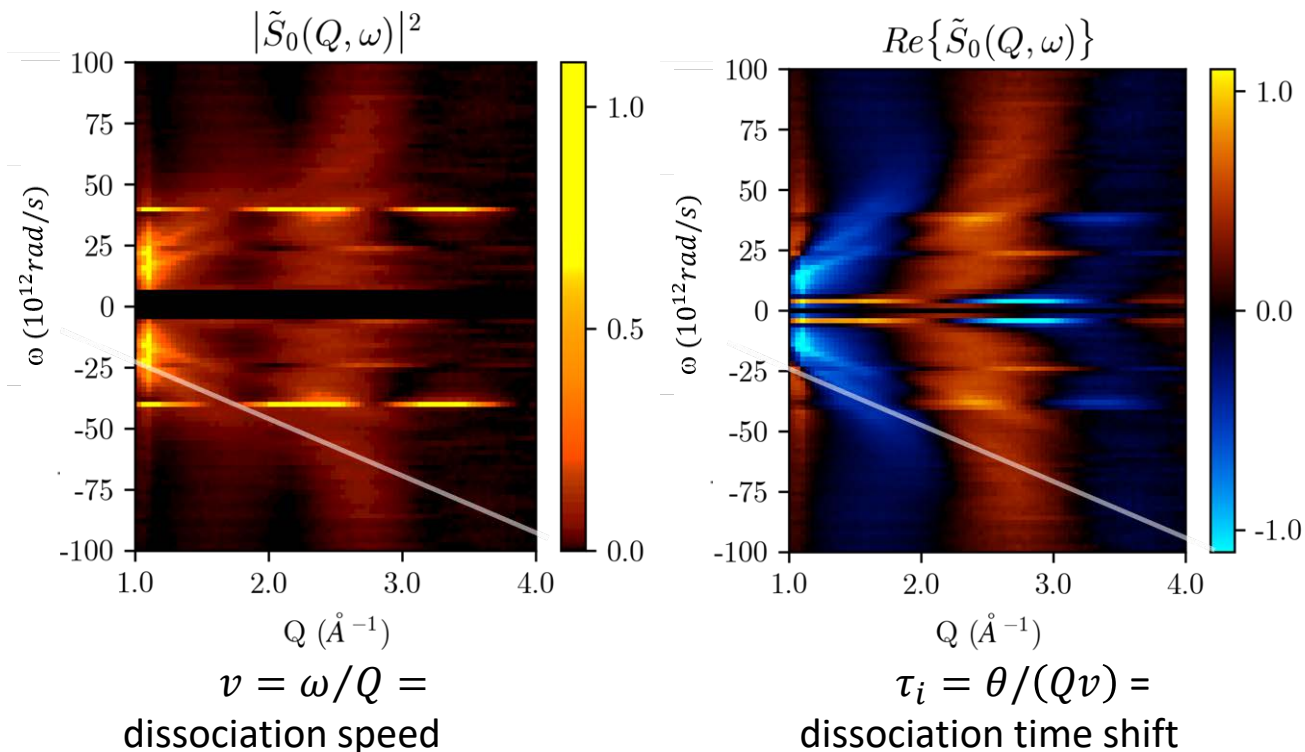
Spontaneous Hyper-Raman Scattering



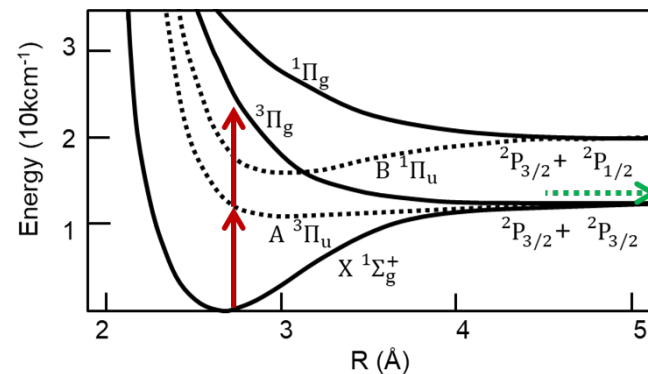
Center of oscillation	$3.10 \pm 0.15 \text{ \AA}$
Oscillation frequency ω	$24 \pm 2 \times 10^{12} \text{ rad/s}$
Oscillation amplitude	$0.4 \pm 0.4 \text{ \AA}$
Oscillation phase δ	0°

P. H. Bucksbaum, M. R. Ware,
A. Natan, J. P. Cryan, and J.
M. Glowia, ArXiv:1911.01323
[Physics] (2019).

Two photon dissociation

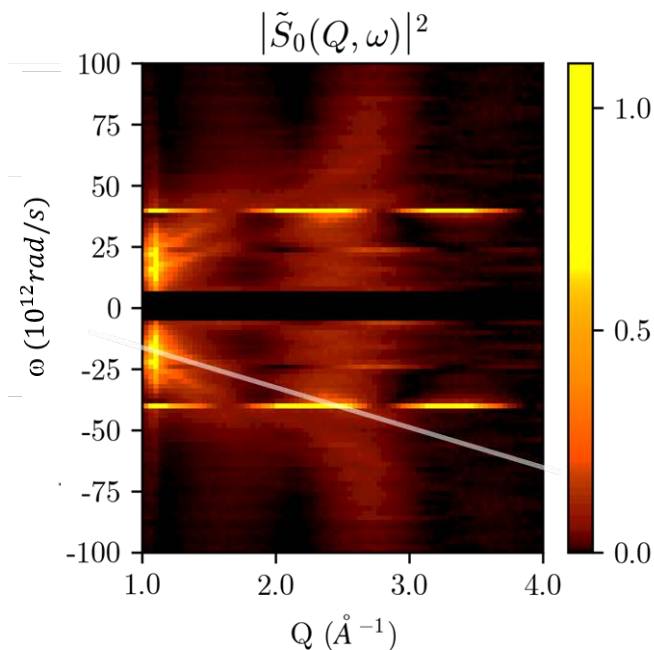


Separation velocity $v_{(d)}$	$23 \pm 1 \text{ \AA/ps}$
Dissociation Time Shift $\tau_{i(d)}$	$+124 \pm 20 \text{ fs}$



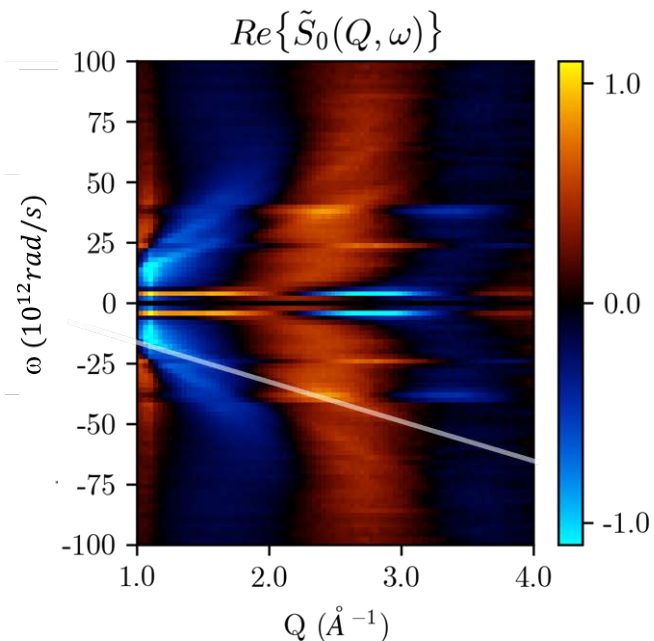
P. H. Bucksbaum, M. R. Ware,
 A. Natan, J. P. Cryan, and J.
 M. Glowia, ArXiv:1911.01323
 [Physics] (2019).

Two photon curve-crossing dissociation



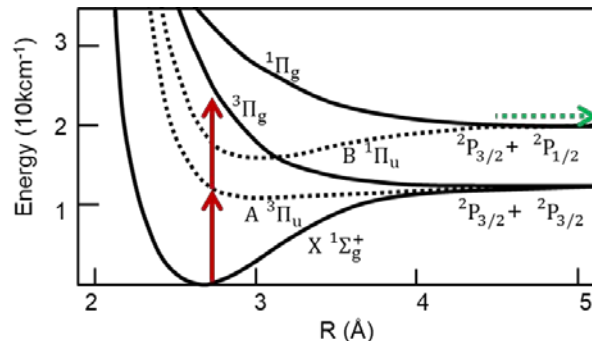
$v = \omega/Q =$
dissociation speed

Separation velocity $v_{(c)}$
Dissociation time shift $\tau_{i(c)}$



$\tau_i = \theta/(Qv) =$
dissociation time shift

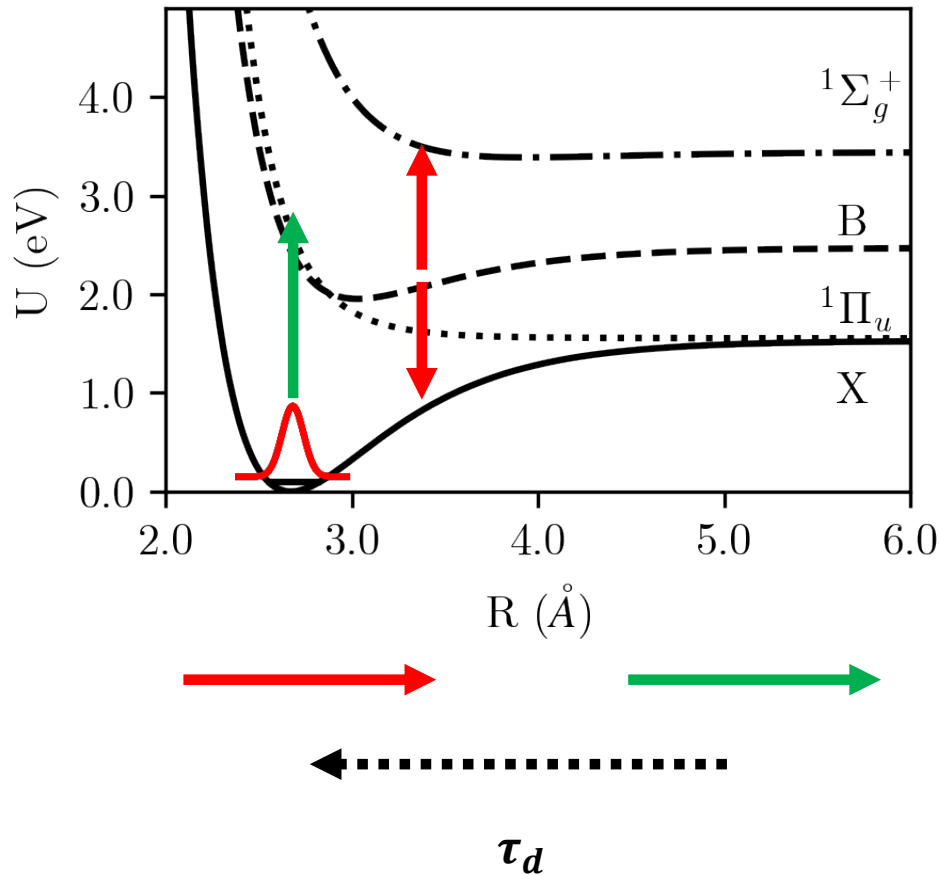
$14 \pm 1 \text{ \AA/ps}$
 $+250 \pm 20 \text{ fs}$



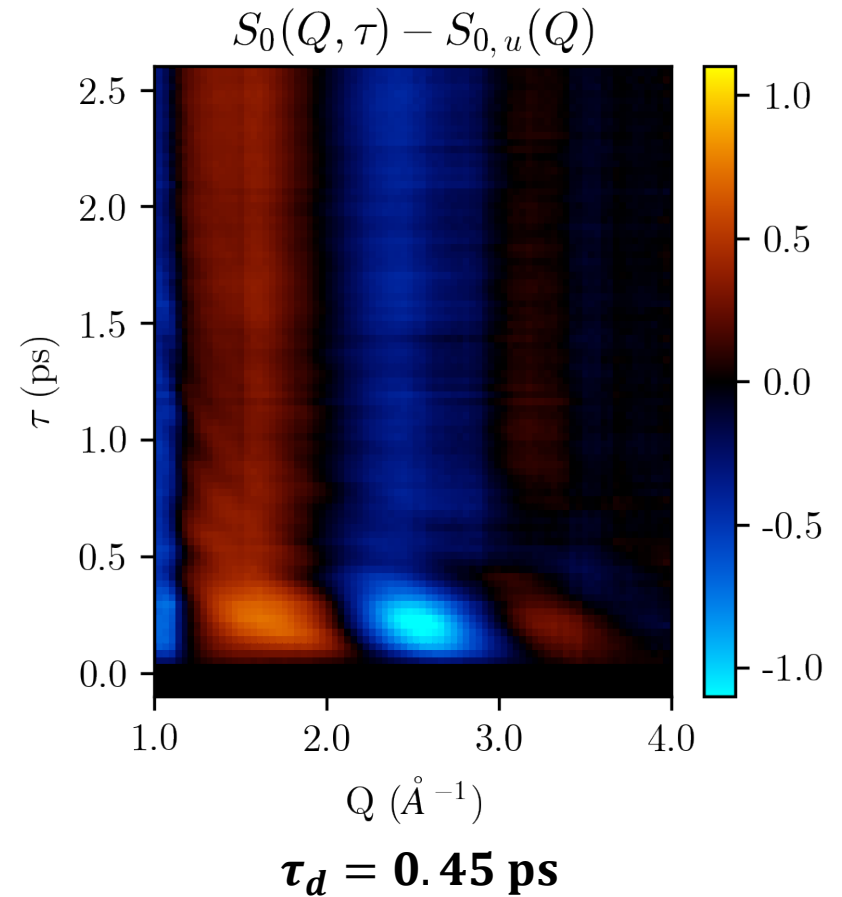
P. H. Bucksbaum, M. R. Ware,
A. Natan, J. P. Cryan, and J.
M. Glowia, ArXiv:1911.01323
[Physics] (2019).

Coherent control: The experiment

Potential energy curves



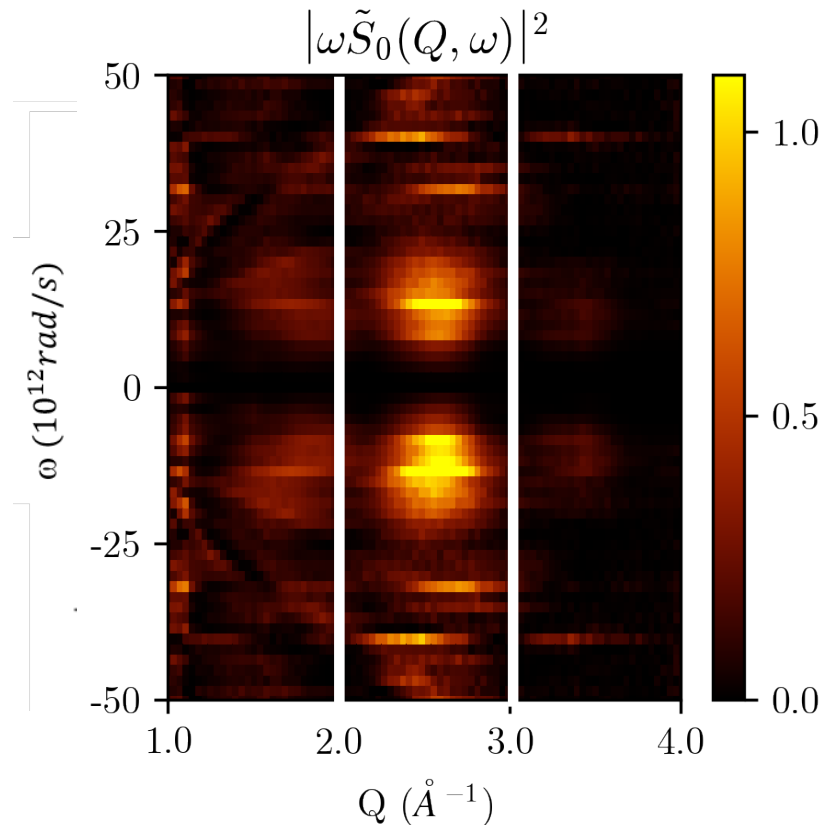
Measurement



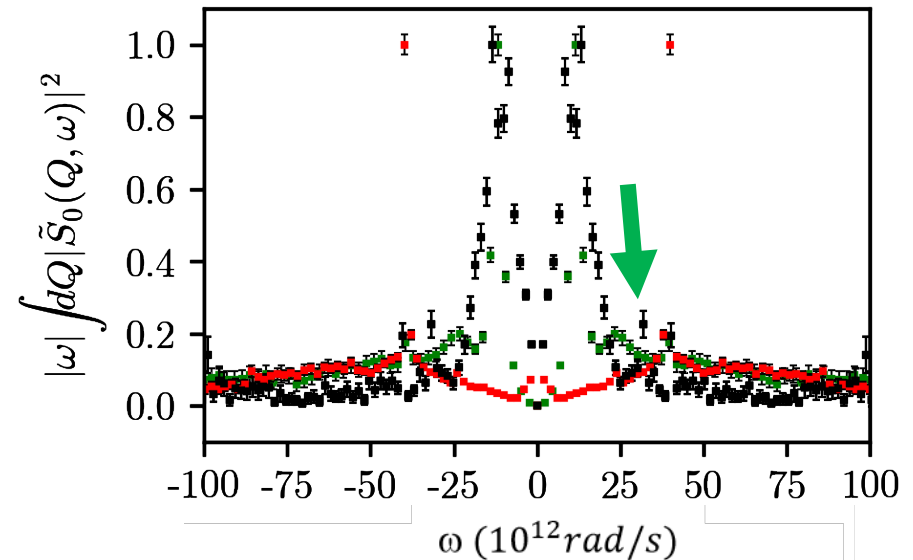
M.R. Ware

Coherent control: Identifying the dumped population

FRXS

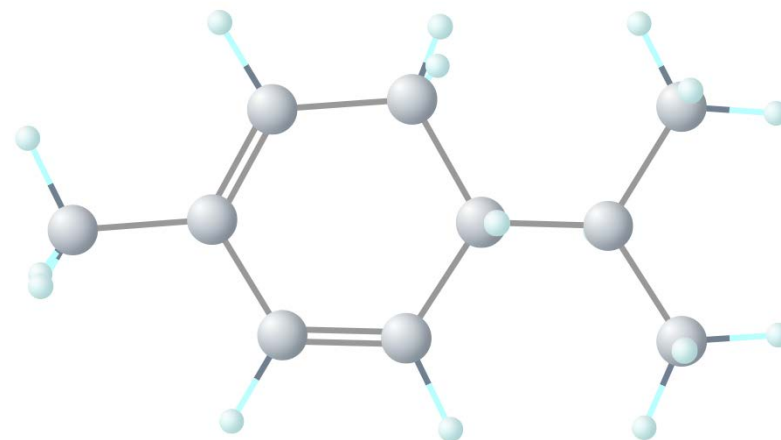
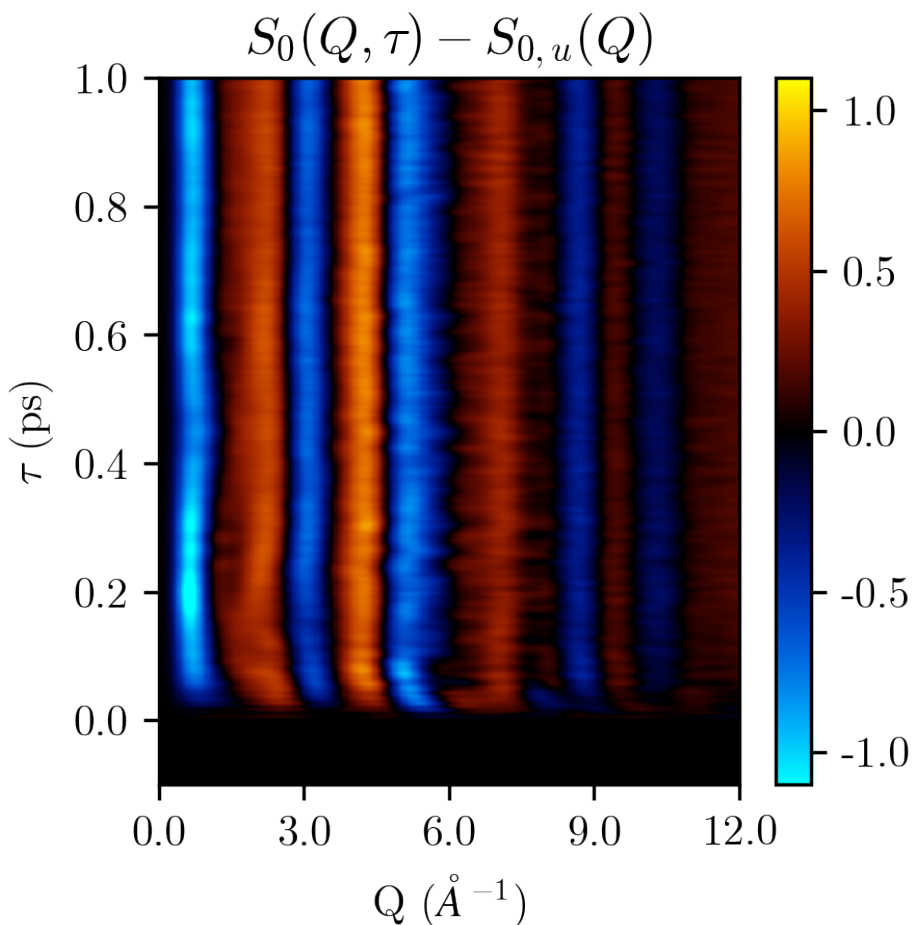


Integrated spectrum



$\omega_d = 31.8 \pm 0.9 \text{ THz}$
as compared to $\omega_{lit} = 31.9 \text{ THz}$

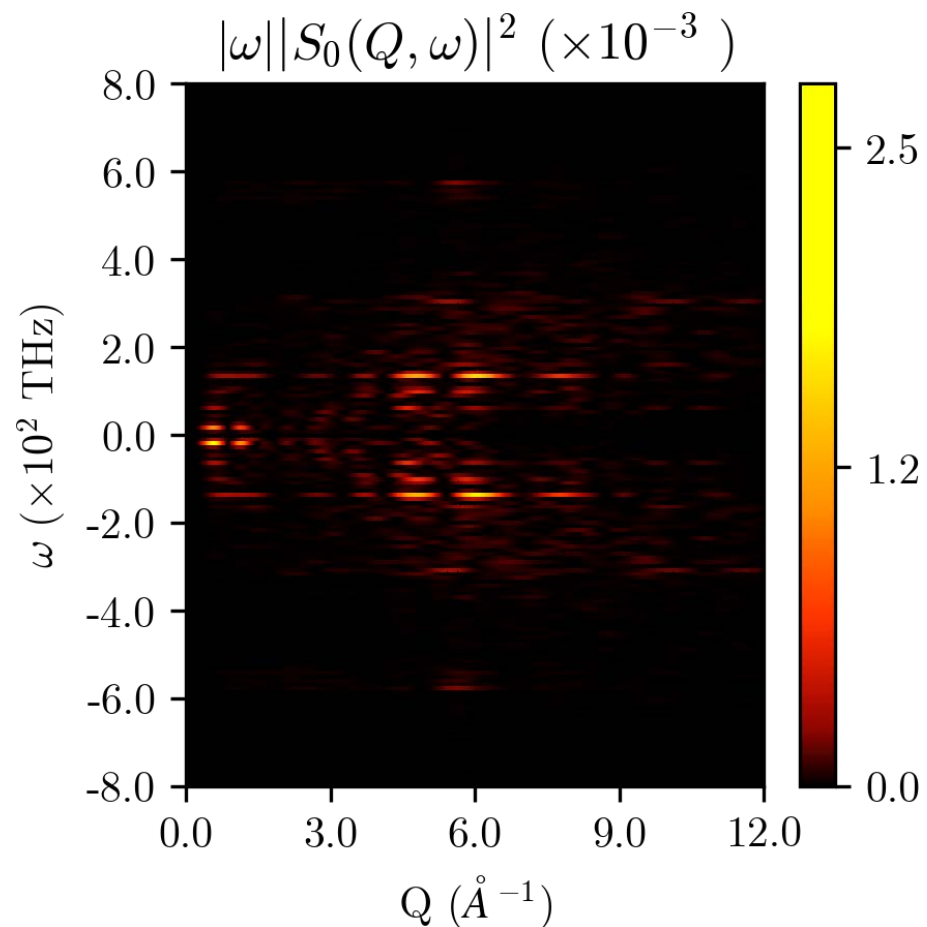
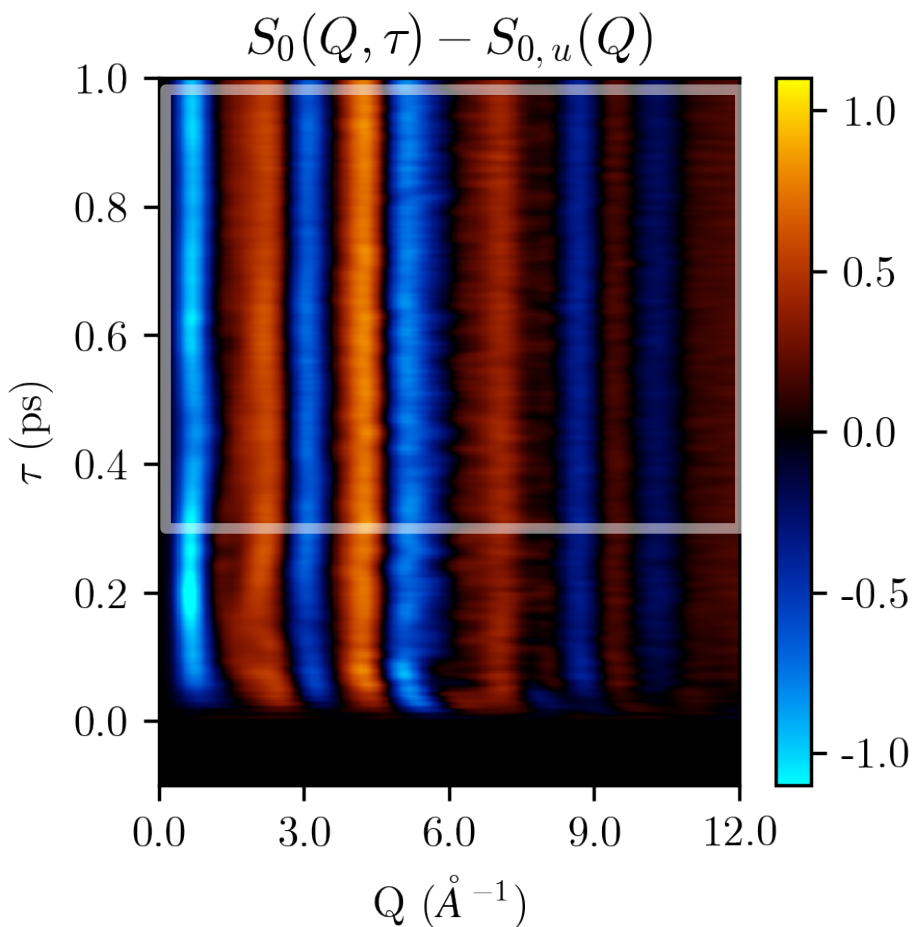
Finding the essential modes of motion in a polyatomic molecule with many modes:



alpha-Phellandrene

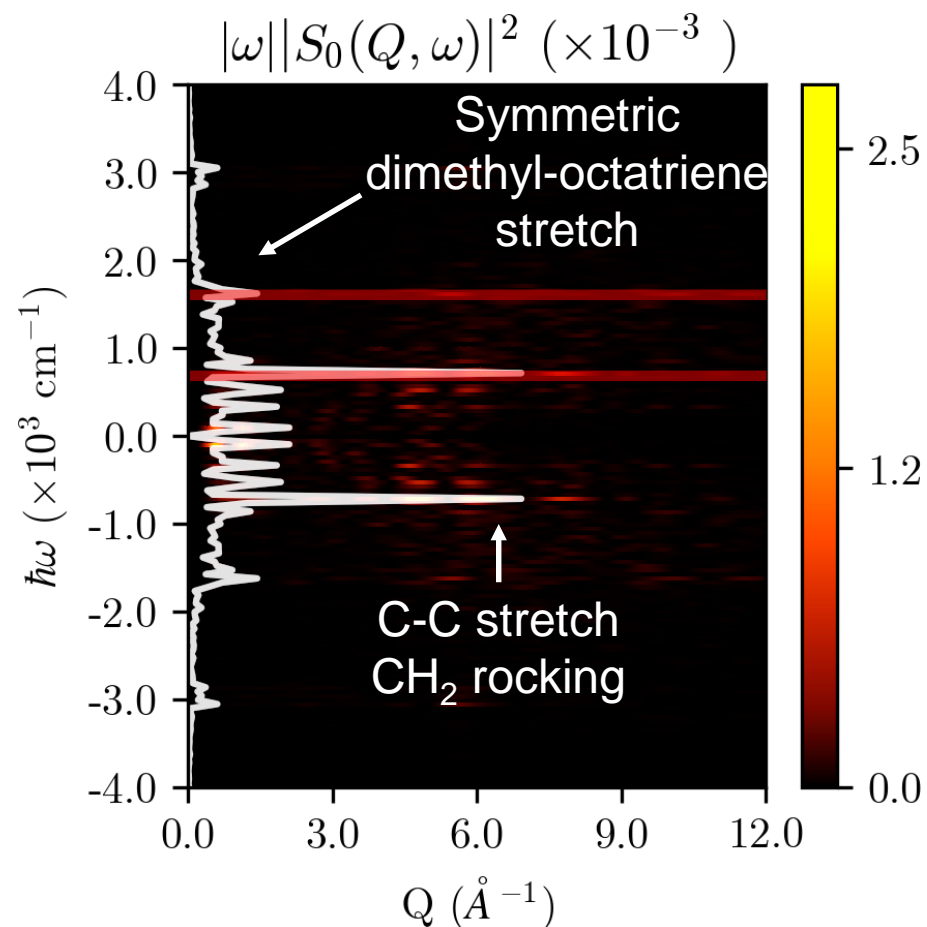
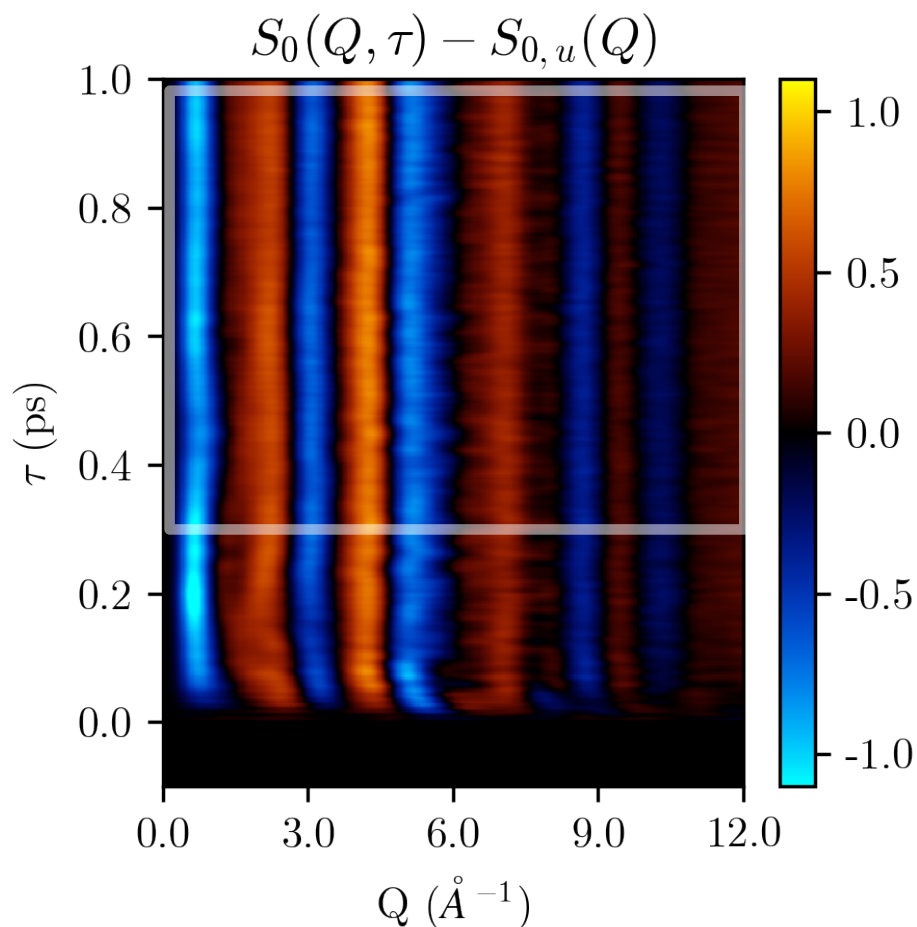
Simulation courtesy of
David Sanchez, Martinez group

Finding the essential modes of motion



Simulation courtesy of
David Sanchez, Martinez group

Finding the essential modes of motion



Simulation courtesy of
David Sanchez, Martinez Group

Reid, P.J, et. al. J. Phys. Chem. 1990.
Reid, P.J, et. al. J. Am. Chem. Soc. 1993.

Conclusion:

- **Optics-based methods cannot measure sub-Angstrom displacement and femtosecond motion in molecules.**
- **Femtosecond X-ray laser scattering fills this gap, and can answer questions about the dynamics of motion:**
 - **How long does it take the atoms to move?**
 - **Where is the motion within the molecule?**
 - **What is the amplitude of this motion?**
 - **What is the phase, i.e. when initially excited, do the molecular bonds first contract or first expand?**
- **Taken together, these are molecular movies that provide new insights about mechanisms in molecular physics and chemistry.**

Credits

- **Adi Natan**
- **Matt Ware**
- **Mike Glownia**
- **James Cryan**
- **Noor Al-Sayyad**
- Ian Gabalski
- Ruaridh Forbes
- Andreas Kaldun
- Anna Wang
- Nick Werby
- Jordan O'Neal
- Andrei Kamalov
- Ryan Coffee
- Robert Hartsock
- David Sanchez
- Miki Minitti
- Todd Martinez
- Rob Parrish
- Joe Subotnik
- David Reis

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